

# Algorithms to compute Stark Numbers

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This presentation is part of an ongoing project that started in the summer of 2022 with Professeur Hugo Chapdelaine (Université Laval).



# Stark Conjecture

## Theorem. Class Number Formula

The function  $\zeta_K(s)$  converges absolutely for  $\Re(s) > 1$  and extends to a meromorphic function defined for all complex  $s$  with only one simple pole at  $s = 1$  with residue

$$\lim_{s \rightarrow 1} (s - 1)\zeta_K(s) = \frac{2^r (2\pi)^c R_K h_K}{\#\mu_K \sqrt{|d_K|}}$$

Using the functional equation we can get some intuition of the first coefficient of the Taylor expansion.

## Rank one Stark Conjecture

Let  $K$  be a number field. For all L function  $L_K(s)$  there exists an algebraic integer  $u \in \mathcal{O}_H$ , where  $H$  is the associated ray class field, such that

$$L'_K(0) = \frac{\delta(v)}{\#\mu_K} \log |\tilde{v}(u)|$$

where  $\delta(v)$  is 1 if  $v$  is real and 2 if it is complex.

Let  $K$  be a Number Field and  $\{v_1, \dots, v_t\}$  the set of places of this field. Given any vector  $p \in (\mathbb{F}_2)^t$ , we define the sign map

$$\omega_p(x) := \prod_{j=1}^t (s_{v_j}(x))^{p_j} \text{ where } s_v(x) = \begin{cases} \text{sign}(v(x)) & \text{if } v \text{ real} \\ 1 & \text{if } v \text{ complex} \end{cases}$$

We can also define the fully positive integers

$$\mathcal{O}_K(\infty) := \{\alpha \in \mathcal{O}_K \setminus \{0\} : s_v(\alpha) = 1 \text{ for all places } v\}$$

Given a lattice  $\mathfrak{n} \subseteq K$

$$\mathcal{V}_{a,b,\mathfrak{n}}^+ := \{\varepsilon \in \mathcal{O}_K^\times(\infty) : (\varepsilon - 1)a \in \mathfrak{n}, (\varepsilon - 1)b \in \mathfrak{n}^*, (\varepsilon - 1)ab \in \mathcal{O}_K^*\}$$

## Definition. Lattice Zeta function

Given a lattice  $\mathfrak{n} \subseteq K$  and  $s \in \mathbb{C}$  such that  $\Re(s) > 1$ , the lattice zeta function of  $\mathfrak{n}$  is defined as

$$Z_{\mathfrak{n}}(a, b, p; s) = \sum'_{x \in (\mathcal{V}_{a,b,\mathfrak{n}}^+ \setminus \mathfrak{n})} \frac{\omega_p(x+a) e^{2\pi i \operatorname{Tr}_K(b(x+a))}}{|\mathbf{N}_K(x+a)|^s}$$

Hugo Chapdelaine proved that this function admits a meromorphic extension to  $\mathbb{C}$ .

# Zeta function associated to a ray class field

Giving a level  $f \in \mathcal{O}_K$  and a fractional ideal  $\mathfrak{a} \subseteq K$ , we define the Zeta function associated to the ray class field of level  $f$  as

$$Z_{f,\mathfrak{a}}(a, b, p; s) := Z_{f\mathfrak{a}^{-1}}(a, b, p; s)$$

## Proposition.

All the representatives of a given class in  $\text{Cl}_f^p(K)$  generate the same zeta function. Consequently, we will compute the Stark Numbers associated to the classes of the ray class group.

## Rank one Stark Conjecture

Let  $\mu_H$  be the set of roots of units in  $H$ . For a given class  $[\mathfrak{a}] \in \text{Cl}_f^p(K)$ , there exists an algebraic number  $u_{[\mathfrak{a}]} \in \mathcal{O}_H$  such that

$$\frac{\partial}{\partial s} Z_{f,\mathfrak{a}}(a, b, p; 0) = -\frac{\delta(\tilde{v}_1)}{\#\mu_H} \log |\tilde{v}_1(u_{[\mathfrak{a}]})|$$

We start assuming  $K$  is a totally real quadratic extension, that  $\tilde{v}_1$  is real and  $Cl^+(K) = 0$ . Chapdelaine proved the following equation for the first derivative of Lattice Zeta functions

$$\frac{\partial}{\partial s} Z_{\mathfrak{a}}(a, b, p; 0) = \left(\frac{\pi}{2}\right)^2 |d_K|^{1/2} Z_{\mathfrak{a}^*}(-b, a, p; 1)$$

Using basic algebra, we have the following expression

$$Z_{(f\mathfrak{a}^{-1})^*}(0, 1, p; 1) = 2i\kappa(\mathbf{N}(f\mathfrak{a}^{-1})) \left( \sum_{\gamma} \frac{\sin(\alpha_{\gamma})}{|\mathbf{N}\gamma|} - \sum_{\gamma} \frac{\sin(\alpha_{\gamma})}{|\mathbf{N}\gamma|} \right)$$

During the Summer of 2022, we computed three decimals at Université Laval for the example  $\mathbb{Q}(\sqrt{29})$ ,  $f = (3 - \sqrt{29})/2$  and  $p = (1, 0)$

4.624

Shintani was able to compute 33 decimals of this number 45 years ago!!!



# Algorithm for the Riemann Zeta Function

## Real analytic Eisenstein series

For  $z \in \mathfrak{h}$ ,  $s \in \mathbb{C}$ , with  $\Re(s) > 1$ , we define the lattice real Eisenstein series as

$$G(z, s) := \sum'_{(m,n) \in \mathbb{Z}^2} \frac{\Im(y)^s}{|mz + n|^{2s}}$$

This function admits a functional equation

$$G(z, s) = \frac{\pi^{-(1-s)} \Gamma(1-s)}{\pi^{-s} \Gamma(s)} G(z, 1-s)$$

The Eisenstein series is 1-periodic and, therefore, admits a Fourier expansion

$$G(x + yi) = a_0(y, s) + \sum_{n \in \mathbb{Z} \setminus \{0\}} a_n(y, s) e^{2\pi n x i}$$

where

$$a_0(y, s) = 2\zeta(2s)y^s + \varphi(s)2\zeta(2s - 1)y^{1-s}$$

$$\varphi(s) := \frac{\pi^{1/2}\Gamma(s - \frac{1}{2})}{\Gamma(s)}$$

and

$$a_n(y, s) = \frac{4|n|^{s-1/2}}{\pi^{-s}\Gamma(s)} \sigma_{1-2s}(|n|) y^{1/2} K_{s-\frac{1}{2}}(2\pi|n|y), \quad \sigma_\mu(|n|) = \sum_{0 < d|n} d^\mu$$

# Colmez's Trick

It is easy to check that

$$G(\gamma z, s) = G(z, s), \quad \gamma \in \mathrm{SL}_2(\mathbb{Z})$$

Defining  $R(z, s) = G(z, s) - a_0(y, s)$ , we have

$$G(z, s) - G(-z^{-1}, s) = 0 \iff$$

$$a_0(y, s) - a_0\left(\frac{y}{x^2 + y^2}, s\right) = -R(z, s) + R(-z^{-1}, s) = F(z, s)$$

# Colmez's Trick

Taking two different elements  $z_1, z_2 \in \mathfrak{h}$ , we get the system of equations

$$2M \begin{pmatrix} \zeta(2s) \\ \varphi(s)\zeta(2s-1) \end{pmatrix} = \begin{pmatrix} F(z_1, s) \\ F(z_2, s) \end{pmatrix}$$

If we specialise to

$$z_1 = \sqrt{a^{-1} - 1} + i, \quad z_2 = \sqrt{a^{-2} - 1} + i, \quad a \in (0, 1)$$

We can give a general expression of the Riemann Zeta function

$$\zeta(2s) = \frac{1}{2} \frac{(1 + a^{1-s})F(z_1, s) - F(z_2, s)}{a^{2s} - a^s + a^{1-s} - a}$$

# Computational Results

A classic example

$$\zeta(2) = 1.64493406684\dots, \quad 9.63 \cdot 10^{-200}$$

A trivial example

$$\zeta(-6) = 0.0000000000\dots, \quad 0.00 \cdot 10^{-200}$$

A non-trivial one

$$\zeta(1/2 + 14.1347251417\dots i) = 1.406844906066\dots \cdot 10^{-200}$$

# Problems that might arise

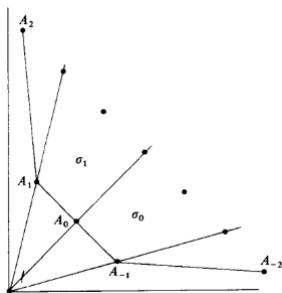
Which problems will we face when we generalize this algorithm to Zeta functions associated to a ray class field?

- 1 The computation of the quotient  $\mathcal{V}_{a,b,f\mathfrak{a}^{-1}}^+ \setminus f\mathfrak{a}^{-1}$
- 2 The generalization of the division algorithm.

For the second problem, simplifying  $Cl^+(K) = 0$  is quite helpful.

$$(\lambda) = l_1 l_2 \cdots l_n$$

# The extremount points



The extremount points satisfy the following equation  $A_{k-1} + A_{k+1} = b_k A_k$  where  $b_k$  are the coefficients of the fraction expansion of  $\tau$ .

# The algorithm for Lattice Zeta functions

## Lattice real analytic Eisenstein series

Given two lattices  $\mathfrak{n}, \mathfrak{m} \subseteq K$ , a vector  $p \in \mathbb{F}_2^2$ ,  $z \in \mathfrak{h}$  and  $U \in M_2(K)$ , we define the following

$$\begin{aligned} G_Q(z, s) &= G_{(\mathfrak{m}, \mathfrak{n})}(U, p; z, s) = \\ &= \sum'_R \frac{\omega_p((m + v_1)z + (n + v_2)) e^{2\pi i \operatorname{Tr}(u_1(m + v_1) + u_2(n + v_2))}}{|\mathbf{N}((m + v_1)z + (n + v_2))|^{2s}} \mathfrak{S}(z)^{1s} \end{aligned}$$

As in the last case, the function is 1-periodic and, therefore, admits a Fourier expansion

$$G_Q(x + y i, s) = a_0(y, s) + \sum_{d \in D \setminus \{0\}} a_d(y, s) e^{2\pi i \operatorname{Tr}(dx)}$$

with the first coefficient equal to

$$e_1 \delta_{\mathfrak{m}}(v_1) Z(v_2, u_2, p; 2s) \mathbf{N}(y)^s + \Omega_p(s) \delta_{\mathfrak{n}^*}(u_2) e_2 Z(v_1, u_2, p; 2s - 1) \mathbf{N}(y)^{1-s}$$



# The algorithm for Lattice Zeta functions

The non-zero coefficients will be

$$a_d(y, s) = \mathbf{N}(y)^s \sqrt{d_K} |\mathbf{N}f| B_d(y, s) b_d(s)$$

with

$$b_d(s) := \sum'_{(\xi_1, \xi_2) \in R_d // \mathcal{O}_K^\times (f_\infty)} (-1)^{W_p(\xi_2)} |\mathbf{N}(\xi_2)|^{2s-1}$$

The equation for the  $S = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \in \mathrm{SL}_2(\mathcal{O}_K)$  action is

$$G_Q(Sz, s) = \omega_{-p}(z) G_{Q^*}(z, s)$$

# Computational results

Chapdelaine worked out an equation that relates the lattice invariants with Shintani's

$$\log X_{f\infty}(\mathfrak{a}) = i \frac{|d_K|^{1/2}}{4\pi} (Z^*(\mathfrak{a}, f\infty, (1, 0); 1) + Z^*(\mathfrak{a}^\sigma, f^\sigma\infty, (1, 0); 1))$$

Let  $K = \mathbb{Q}(\sqrt{29})$ ,  $f = (3 - \sqrt{29})/2$  and  $p = (1, 0)$ . We use magma to compute the ray class group in this case

$$\text{Cl}_f^p(K) = \{[\mathcal{O}_K], [2\mathcal{O}_K]\} \cong \mathbb{Z}/2\mathbb{Z}$$

Using the algorithm we compute all the values present in the equation

$$Z^*(\mathcal{O}_K, f\infty, (1, 0); 1) = -1.178 \dots 10^{-27} + 3.573 \dots i$$

$$Z^*(\mathcal{O}_K^\sigma, f^\sigma\infty, (1, 0); 1) = -2.984 \dots 10^{-28} + 1.051 \dots 10^{-27} i$$

$$Z^*(2\mathcal{O}_K, f\infty, (1, 0); 1) = 8.469 \dots 10^{-28} + 1.365 \dots i$$

$$Z^*((2\mathcal{O}_K)^\sigma, f^\sigma\infty, (1, 0); 1) = 7.688 \dots 10^{-28} - 3.495 \dots 10^{-27} i$$

# Computational results

With these values, we compute Shintani's invariant up to an error of  $10^{-54}$

$$X_f(\mathcal{O}_K) = 4.6242866227 \dots, \quad X_f(2\mathcal{O}_K) = 1.7949175568 \dots$$

Using the LLL-algorithm we get the expected polynomial for these numbers

$$x^8 - 9x^7 + 29x^6 - 52x^5 + 63x^4 - 52x^3 + 29x^2 - 9x + 1$$

Magma even give us the decomposition of such polynomial over  $\mathcal{O}_F[x]$

$$\left( x^4 - \frac{9 + \sqrt{29}}{2}x^3 + (8 + \sqrt{29})x^2 - \frac{9 + \sqrt{29}}{2}x + 1 \right) \cdot$$
$$\left( x^4 + \frac{-9 + \sqrt{29}}{2}x^3 + (8 - \sqrt{29})x^2 + \frac{-9 + \sqrt{29}}{2}x + 1 \right)$$

# Charollois-Darmon Conjecture

Charollois and Darmon consider the L function

$$L(M, s) := (\mathbf{NM})^s \delta_l \sum'_{x \in M/V} \frac{\text{sign}(\mathbf{N}(x))}{|\mathbf{N}(x)|^s}$$

and the associated Eisenstein series

$$E(z_1, \dots, z_n) = \zeta_F(-1) + \sum_{\mu \in \mathcal{O}_F, \mu \gg 0} \left( \sum_{(\nu) | (\mu)} |\mathbf{N}(\nu)| \right) e^{2i\pi \left( \frac{\mu_1}{\sigma_1} z_1 + \dots + \frac{\mu_n}{\sigma_n} z_n \right)}$$

Introducing the notation  $\omega_E = E(z_1, \dots, z_n) dz_1 \wedge \dots \wedge dz_n$ , we can define the  $n$ -differential form

$$\omega_{\text{Eis}} := \begin{cases} \frac{(2\pi i)^n}{\sqrt{d_F}} \omega_E + \frac{R_F}{2} \left( \frac{dz_1 \wedge d\bar{z}_1}{y_1^2} - \frac{dz_2 \wedge d\bar{z}_2}{y_2^2} \right) & \text{if } n = 2 \\ \frac{(2\pi i)^n}{\sqrt{d_F}} \omega_E & \text{if } n > 2 \end{cases}$$

# Charollois-Darmon Conjecture

The form  $\omega_{\text{Eis}}^+ = \frac{1}{2}(\text{Id} + T_v^*)\omega_{\text{Eis}}$  is "the complex real part" of the form  $\omega_{\text{Eis}}$ .

## Theorem (Charollois-Darmon, 2008)

Let  $M$  be an  $\mathcal{O}_I$ -module associated to a  $\tau \in K$ . The first derivative of the L function satisfies

$$(-2)^{n-1} \int_{C_\tau} \omega_{\text{Eis}}^+ = (2i\pi)^{n-1} L'(M, 0)$$

## Charollois-Darmon Conjecture (2008)

For all  $\mathcal{O}_I$ -module  $M$  associated to  $\tau \in K$ , there exists an algebraic number  $u_\tau \in \mathcal{O}_H$  such that

$$e_I \frac{(-2)^{n-1}}{n_F} \int_{\partial C = n_F \Delta_\tau} \omega_{\text{Eis}} = 2\delta_I (2i\pi)^{n-1} \log(\tilde{v}_1(u_\tau))$$

# Generalization Charollois-Darmon Conjecture

For the generalization to lattice zeta functions, we define the lattice Eisenstein series of weight two

$$E_2^*(U, Z) = \sum_{(m,n) \in \left(\mathcal{O}_F^2 / \mathcal{V}_{0,b,\mathcal{O}_F}^+\right)} \frac{e^{-2\pi i \operatorname{Tr}(u_1(\delta m + v_1) + u_2(\delta n + v_2))}}{\mathbf{N}((\delta m + v_1)z + (\delta n + v_2))^2}$$

Using the same notation as the last case

$$\Omega_{E_2^*}(b; z) := E_2^*(b; z) dz_1 \wedge \cdots \wedge dz_n$$

we define the  $n$ -differential form

$$\Omega_{Eis}(b; z) := \begin{cases} -2\Omega_{E_2^*}(b; z) + \frac{R_F}{2} \left( \frac{dz_1 \wedge d\bar{z}_1}{y_1^2} - \frac{dz_2 \wedge d\bar{z}_2}{y_2^2} \right) & \text{if } n = 2 \\ (-2i)^n \Omega_{E_2^*}(b; z) & \text{if } n > 2 \end{cases}$$

# Generalization Charollois-Darmon Conjecture

Similarly to the last case the form  $\Omega_{\text{Eis}}^+ = \frac{1}{2}(U_{v_j} + T_{v_j})\Omega_{\text{Eis}}$  emulates the real part of  $\Omega_{\text{Eis}}$ .

## Theorem (Caralps, 2023)

Given a class  $[\mathfrak{a}] \in \text{Cl}_f^+(K)$  associated to a  $\tau \in K$ , the following equation is satisfied

$$\int_{C_\tau} \Omega_{\text{Eis}}^+ = \left( \frac{\mathbf{N}_{K/\mathbb{Q}}(\nu)\sqrt{\pi}}{d_f\sqrt{d_F}} \right) \frac{\partial}{\partial s} Z_{[\mathfrak{a}]}^\omega(b; 0)$$

## Generalization Charollois-Darmon Conjecture (Caralps, 2023)

For all  $\mathcal{O}_K$ -module  $\mathfrak{n}$  with associated elements  $\nu \in \mathcal{O}_F(\infty)$  and  $\tau \in K$ , that satisfy  $[\nu(\mathcal{O}_F + \tau\mathcal{O}_F)] = [(\mathfrak{n}^{-1})^*]$ , the Stark number  $u_\tau \in \mathcal{O}_H$  of  $Z_{[\mathfrak{n}]}^\omega(b; s)$  satisfies the equation

$$\frac{\#\mu_H}{n_F} \int_{\partial C = n_F \Delta_\tau} \Omega_{\text{Eis}} = -2 \left( \frac{\mathbf{N}_{K/\mathbb{Q}}(\nu)\sqrt{\pi}}{d_F\sqrt{d_F}} \right) \log(\tilde{v}_1(u_\tau))$$

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